

# Fluctuations of the Lyapunov exponents in chaotic extended systems

Diego Pazó<sup>1</sup>, Juan M. López<sup>1</sup> and Antonio Politi<sup>2</sup>

<sup>1</sup>Instituto de Física de Cantabria (IFCA), CSIC-Universidad de Cantabria, 39005 Santander, Spain

<sup>2</sup>Institute for Complex Systems and Mathematical Biology and SUPA, University of Aberdeen, Aberdeen AB24 3UE, United Kingdom

Finite-time Lyapunov exponents of generic chaotic dynamical systems fluctuate in time. These fluctuations are due to the different degree of stability across the accessible phase-space. A numerical study [1] has revealed that the diffusion coefficient  $D$  of the Lyapunov exponents (LEs) exhibits a non-trivial scaling behavior,  $D(L) \sim L^{-\gamma}$ , with the system size  $L$ .

For chaotic dissipative systems, we show that the wandering exponent  $\gamma$  can be expressed in terms of the roughening exponents associated with the corresponding “Lyapunov-surface”. Our theoretical predictions are supported by the numerical analysis of several spatially-extended systems [2]. In particular, we find that the wandering exponent of the first LE is universal: in view of the known relationship with the Kardar-Parisi-Zhang equation,  $\gamma$  can be expressed in terms of known critical exponents (see typical example in Fig. 1). Furthermore, our simulations reveal that the bulk of the spectrum exhibits a clearly different behavior and suggest that it belongs to a possibly unique universality class, which has, however, yet to be identified.

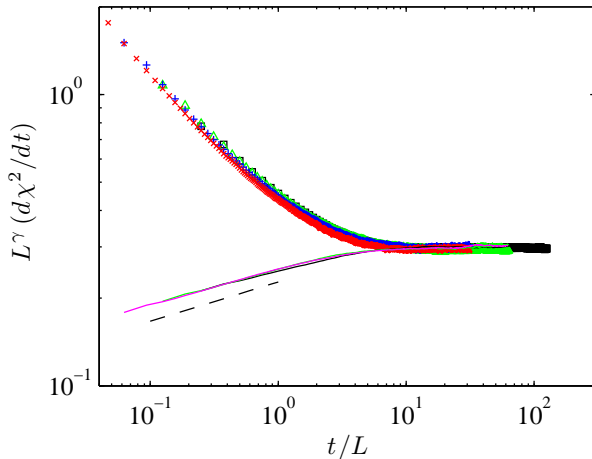


Figure 1: Scaled FTLE fluctuations in a chain of coupled Hénon maps. The exponent  $\gamma$  is set equal to  $\gamma = 0.865$ . The solid curves (which correspond to  $L = 40, 80,$  and  $160$ ) have been obtained for the 0-norm. The various symbols (squares, triangles, pluses, and crosses correspond to  $L = 40, 80, 160$  and  $320$ , respectively) are obtained by using the Euclidean norm of the Gram-Schmidt LVs. The dashed line corresponds to a power law growth  $(t/L^z)^\nu$  with  $\nu = 0.135$  and  $z = 1$ .

In contrast, in generic Hamiltonian lattices (see example in Fig. 2) the diffusion coefficient of the maximum Lyapunov exponent diverges in the thermodynamic limit [3]. We trace this back to the long-range correlations associ-

ated with the evolution of the hydrodynamic modes. In the case of normal heat transport, the divergence is even stronger, leading to the breakdown of the usual single-function Family-Vicsek scaling ansatz. A similar scenario is expected to arise in the evolution of rough interfaces in the presence of a suitably correlated background noise.

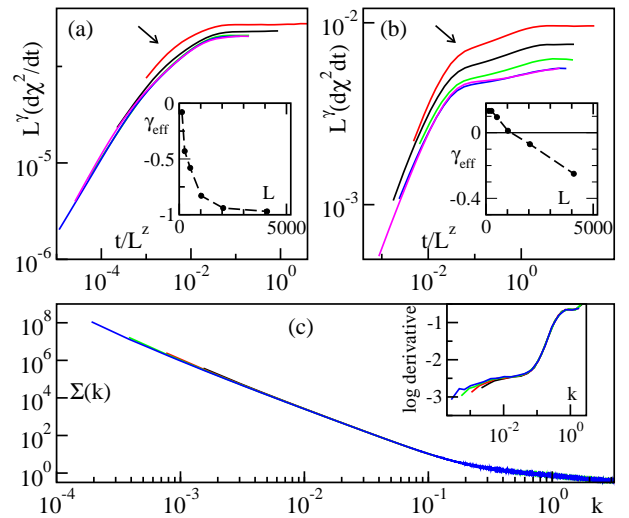


Figure 2: Scaled FTLE fluctuations for the  $\Phi^4$  (panel a) and FPU (panel b) models, according to the theoretical ansatz, for sizes  $L = 256, 512, 1024, 2048, 4096$ . The derivative  $d\chi^2/dt$  is considered because of the faster temporal convergence to the asymptotic results. The optimal collapse of the data sets for the larger system sizes is achieved setting  $\gamma = -0.97$  and  $z = 2.15$  in (a) and  $\gamma = -0.25$  and  $z = 1.45$  in (b). The insets show the convergence of  $\gamma$ , determined by comparing the sizes  $L$  and  $L/2$ . In panel (c) the structure factor of Lyapunov-vector surface the FPU model is plotted for  $L = 4096, 8192, 16384,$  and  $32768$  (see black, red, green, and blue lines). The logarithmic derivative is plotted in the inset, adopting the same color coding.

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